



## Solving nonlinear fractional autonomous dynamical systems using a Laplace-Daftardar-Jafari method

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### Abstract

In this manuscript, we present a novel analytical method, namely the Laplace Daftardar Jafari method (LDJM), for solving nonlinear fractional-order autonomous dynamical systems using the Caputo fractional derivative. The proposed method is demonstrated to be highly accurate and efficient through various examples, graphical illustrations, and numerical comparisons with the Laplace residual power series method (LRPS). The LDJM has the potential to be applied to a wide range of complex dynamical systems, including biological and physical systems.

**Keywords:** Nonlinear fractional-order systems, Autonomous dynamical systems, Laplace Daftardar-Jafari method, Caputo fractional derivative.


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### 1. Introduction

Autonomous nonlinear dynamical systems are complex systems that appear in various scientific and engineering fields, such as physics, engineering, and economics [28, 6, 10]. These systems are characterized by the absence of external energy sources, and therefore their behavior is solely determined by the initial conditions and the governing laws. Therefore, the fractional derivatives are particularly suitable for such cases, as they can capture the long-term effects of these systems [11, 7, 29]. As a result, numerous researchers investigate these systems to understand nonlinear behavior, analyze stability and instability, develop strategies and methods for controlling their behavior, and comprehend the complex phenomena that arise in them [8, 17, 12, 18, 32, 14]. However, solving these equations can be difficult, especially when they are nonlinear. To resolve this challenge, numerous analytical and numerical methods have been developed to solve differential equations; for instance, the Adomian decomposition method (ADM) [23], the Hussein–Jassim method (HJM) [25], the Chebyshev spectral collocation method [37], the Sawi decomposition method [22], the integral transforms methods [5], the Sumudu decomposition method (SDM) [15], the Elzaki homotopy analysis method [34], the modification homotopy perturbation method [13], the Laplace variational iteration method [24], the Sumudu variational iteration method [27], and other

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methods [26, 36, 2, 19, 1, 20, 21, 33]. Among these methods, the Daftardar-Jafari method (DJM) is one of the elective methods developed in 2006 [16]; it is a modified version of Adomians decomposition method (ADM) [3], which eliminates the need for calculating Adomian polynomials in any of the iterations. The DJM has been extensively studied by many researchers; see [30, 35, 4, 16].

This study aims to solve autonomous nonlinear dynamical systems of fractional order in the Caputo sense, using a novel hybrid approach that reduces computational time, effort, and error. This hybrid method combines the Laplace transform method (LTM) with the Daftardar-Jafari method (DJM), known as the Laplace-Daftardar-Jafari method (LDJM). Furthermore, numerical results will be compared with those obtained using the Laplace residual power series method (LRPS) to provide a comprehensive and effective analysis for solving these complex systems. This comparative analysis will highlight the advantages and limitations of each method, thereby offering insights into the most efficient techniques for tackling autonomous nonlinear dynamical systems of fractional order. By demonstrating the effectiveness of the LDJM, we aim to enhance the understanding and application of fractional calculus in real-world scenarios.

This paper is organized as follows. In Section 2, we present some key definitions of fractional calculus. In Section 3, the LDJM is scrutinized. Section 4 provides illustrative applications of the proposed method. The paper concludes with section 5.

## 2. Preliminaries and Basic Definitions

This section provides an overview of the definitions of fractional calculus.

**Definition 2.1.** The Laplace transform of a function  $y(t)$  is defined as: [7, 20]

$$L[y(t)] = \int_0^{\infty} y(t)e^{-pt} dt = Y(p), \quad t \geq 0 \quad (2.1)$$

**Definition 2.2.** The Caputo fractional derivative of a function  $y(t)$  is defined as: [7, 20]

$$D_t^\gamma y(t) = \frac{1}{\Gamma(m-\gamma)} \int_0^t \frac{y^{(m)}(x)}{(t-x)^{\gamma-m+1}} dx, \quad (2.2)$$

for  $m-1 < \gamma < m$ ,  $m \in \mathbb{N}$ .

**Definition 2.3.** The Laplace transform of the Caputo fractional derivative is defined as:[7, 9]

$$L[D_t^\gamma y(t)] = p^\gamma Y(p) - \sum_{i=0}^{m-1} p^{\gamma-i-1} y^{(i)}(0), \quad (2.3)$$

for  $m-1 < \gamma < m$ ,  $m \in \mathbb{N}$ .

## 3. Analysis of the Method

Consider the following fractional differential equation:

$$D_t^\gamma y(t) = \mathcal{R}[y(t)] + \mathcal{M}[y(t)], \quad 0 < \gamma < 1, \quad t > 0 \quad (3.1)$$

with initial condition

$$y(0) = \varphi, \quad (3.2)$$

where  $D_t^\gamma$  is the Caputo fractional derivative,  $\mathcal{R}$  and  $\mathcal{M}$  are linear and nonlinear operators acting on  $y(t)$  respectively.

Applying the Laplace transform to both sides of equation (4) and using the initial condition (5), we obtain

$$y(t) = \varphi + \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [\mathcal{R}[y(t)] + \mathcal{M}[y(t)]] \right]. \quad (3.3)$$

### 3.1. The Adomian Decomposition Method (ADM)

Suppose the solution of equation (6) takes the form

$$y(t) = \sum_{i=0}^{\infty} y_i(t), \quad (3.4)$$

and the nonlinear function  $M[y(t)]$  is formulated as

$$M \left[ \sum_{i=0}^{\infty} y_i(t) \right] = \sum_{i=0}^{\infty} A_i(t), \quad (3.5)$$

where the Adomian polynomials  $A_i(t)$  are obtained through the relation

$$A_i(t) = \frac{1}{i!} \frac{d^i}{d\lambda^i} M \left( \sum_{n=0}^i \lambda^n y_n(t) \right) \Big|_{\lambda=0}, \quad i \geq 0. \quad (3.6)$$

As a result, the terms of the series are given by:

$$\begin{cases} y_0(t) = \varphi, \\ y_i(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [\mathcal{R}[y_{i-1}(t)] + A_{i-1}(t)] \right], \quad i \geq 1. \end{cases} \quad (3.7)$$

Refer to [14] for a discussion on the convergence of this series.

### 3.2. The Daftardar-Jafari Method (DJM)

In recent years, Daftardar-Gejji and Jafari [24] employed a different decomposition approach for the nonlinear term  $M[y(t)]$ , decomposing it as follows:

$$M \left[ \sum_{i=0}^{\infty} y_i(t) \right] = M[y_0(t)] + \sum_{i=1}^{\infty} \left( M \left[ \sum_{n=0}^i y_n(t) \right] - M \left[ \sum_{n=0}^{i-1} y_n(t) \right] \right). \quad (3.8)$$

The recurrence formula is formulated as:

$$\begin{cases} y_0(t) = \varphi, \\ y_1(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [\mathcal{R}[y_0(t)] + M[y_0(t)]] \right], \\ y_i(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [\mathcal{R}[y_{i-1}(t)] + M \left[ \sum_{n=0}^{i-1} y_n(t) \right] - M \left[ \sum_{n=0}^{i-2} y_n(t) \right]] \right], \quad i \geq 2. \end{cases} \quad (3.9)$$

## 4. Application and Numerical results

In this section, we illustrate two examples of fractional nonlinear autonomous systems of varying orders. The implementation steps of the LDJM are elucidated, and the effectiveness of the proposed approach is validated through a comprehensive analysis, including graphical representations and tabular data.

**Example 4.1.** Firstly, consider the following autonomous 2-dimensional fractional nonlinear systems [1]:

$$\begin{cases} D_t^\gamma x(t) = 0.5x(t), \\ D_t^\gamma y(t) = x^2(t) + y(t), \end{cases} \quad 0 < \gamma < 1, \quad (4.1)$$

subject to initial conditions

$$x(0) = 1, \quad y(0) = 0. \tag{4.2}$$

Operating Laplace transform on both sides of (13) and using (14), we get

$$\begin{cases} x(t) = 1 + 0.5\mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [x(t)] \right], \\ y(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} [\sum_{i=0}^\infty A_i(t) + y(t)] \right], \end{cases} \tag{4.3}$$

where  $A_i(t)$  denote the Daftardar-Jafari polynomials corresponding to the nonlinear operator  $x^2(t)$ .

In view of LDJM,

$$\begin{aligned} x_0(t) &= 1, \quad y_0(t) = 0, \\ x_1(t) &= \frac{0.5t^\gamma}{\Gamma(\gamma + 1)}, \quad y_1(t) = \frac{t^\gamma}{\Gamma(\gamma + 1)}, \\ x_2(t) &= \frac{0.25t^{2\gamma}}{\Gamma(2\gamma + 1)}, \quad y_2(t) = \frac{2t^{2\gamma}}{\Gamma(2\gamma + 1)} + \frac{0.25\Gamma(2\gamma + 1)}{\Gamma^2(\gamma + 1)\Gamma(3\gamma + 1)}t^{3\gamma}, \\ x_3(t) &= \frac{0.125t^{3\gamma}}{\Gamma(3\gamma + 1)}, \\ y_3(t) &= \frac{2.5}{\Gamma(3\gamma + 1)}t^{3\gamma} + \frac{0.25\Gamma(2\gamma + 1)}{\Gamma^2(\gamma + 1)\Gamma(4\gamma + 1)} \left[ 1 + \frac{\Gamma(\gamma + 1)\Gamma(3\gamma + 1)}{\Gamma^2(2\gamma + 1)} \right] t^{4\gamma} \\ &\quad + \frac{0.0625\Gamma(4\gamma + 1)}{\Gamma^2(2\gamma + 1)\Gamma(5\gamma + 1)}t^{5\gamma}. \end{aligned}$$

Thus, the approximate solution of system (13) is

$$\begin{aligned} x(t) &= 1 + \frac{0.5t^\gamma}{\Gamma(\gamma + 1)} + \frac{0.25t^{2\gamma}}{\Gamma(2\gamma + 1)} + \frac{0.125t^{3\gamma}}{\Gamma(3\gamma + 1)} + \dots, \\ y(t) &= \frac{t^\gamma}{\Gamma(\gamma + 1)} + \frac{2t^{2\gamma}}{\Gamma(2\gamma + 1)} + \frac{0.25\Gamma(2\gamma + 1)}{\Gamma^2(\gamma + 1)\Gamma(3\gamma + 1)}t^{3\gamma} + \frac{2.5}{\Gamma(3\gamma + 1)}t^{3\gamma} \\ &\quad + \frac{0.25\Gamma(2\gamma + 1)}{\Gamma^2(\gamma + 1)\Gamma(4\gamma + 1)} \left[ 1 + \frac{\Gamma(\gamma + 1)\Gamma(3\gamma + 1)}{\Gamma^2(2\gamma + 1)} \right] t^{4\gamma} \\ &\quad + \frac{0.0625\Gamma(4\gamma + 1)}{\Gamma^2(2\gamma + 1)\Gamma(5\gamma + 1)}t^{5\gamma} + \dots. \end{aligned}$$

In particular, if we put  $\gamma = 1$ , we have the exact solution in closed form

$$\begin{cases} x(t) = 1 + \frac{1}{2}t + \frac{(\frac{1}{2}t)^2}{2!} + \frac{(\frac{1}{2}t)^3}{3!} + \dots \cong e^{\frac{1}{2}t}, \\ y(t) = t + t^2 + \frac{1}{12}t^3 + \frac{5}{128}t^4 + \frac{1}{320}t^5 + \dots \cong te^t. \end{cases} \tag{4.4}$$

Table 2: A Comparison of absolute errors obtained for example 1 using LDJM and LRPS when  $\gamma = 1$ .

t	LDJM solution		LRPS solution	
	x(t)	y(t)	x(t)	y(t)
0	0.000000	0.000000	0.000000	0.000000
0.1	0.000000	0.000012	0.000000	0.000017
0.2	0.000004	0.000196	0.000004	0.000281
0.3	0.000022	0.001028	0.000022	0.001458
0.4	0.000069	0.003365	0.000069	0.004730
0.5	0.000171	0.008508	0.000171	0.011861

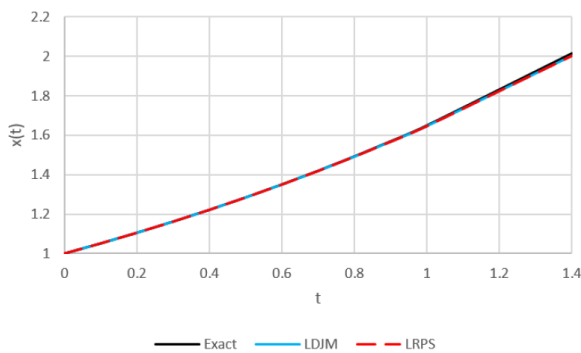
Table 1: A Comparison of numerical results obtained for example 1 using LDJM and LRPS when  $\gamma = 1$ .

t	Exact solution		LDJM solution		LRPS solution	
	x(t)	y(t)	x(t)	y(t)	x(t)	y(t)
0	1.000000	0.000000	1.000000	0.000000	1.000000	0.000000
0.1	1.051271	0.110517	1.051271	0.110505	1.051271	0.110500
0.2	1.105171	0.244281	1.105167	0.244084	1.105167	0.244000
0.3	1.161834	0.404958	1.161813	0.403929	1.161813	0.403500
0.4	1.221403	0.596730	1.221333	0.593365	1.221333	0.592000
0.5	1.284025	0.824361	1.283854	0.815853	1.283854	0.812500

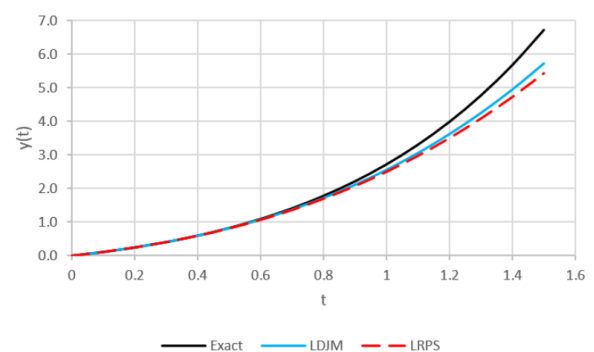
Tab .1 presents the numerical results for the approximate solution of example 1 (4-term solution) at  $\gamma = 1$  and compares them with the exact solution and LRPS method [1]. Tab .2 shows a comparison of absolute errors. Furthermore, Tab .3 displays the approximate solutions obtained by LDJM, for Example .4.1 at different values of  $\gamma$ . Fig. 1 illustrates the comparison between the LDJM and LRPS methods for  $x(t)$  and  $y(t)$  at  $\gamma = 1$ .

Table 3: Numerical results obtained for example 1 using LDJM for different value of  $\gamma$ .

t	$\gamma = 0.9$		$\gamma = 0.7$		$\gamma = 0.5$	
	x(t)	y(t)	x(t)	y(t)	x(t)	y(t)
0	1.000000	0.000000	1.000000	0.000000	1.000000	0.000000
0.1	1.066818	0.142279	1.111243	0.150740	1.179709	0.367889
0.2	1.127716	0.292428	1.184315	0.310995	1.257689	0.553165
0.3	1.188725	0.466284	1.250654	0.498776	1.321539	0.736956
0.4	1.251105	0.670329	1.314838	0.722049	1.379834	0.942428
0.5	1.315519	0.910049	1.378931	0.987577	1.436069	1.182673



(a)



(b)

Figure 1: A comparison of  $x(t)$  and  $y(t)$  solutions between exact, LDJM, and LRPS methods for Example 4.1 when  $\gamma = 1$ .

**Example 4.2.** Secondly, consider the following autonomous 3-dimensional fractional nonlinear systems [1]:

$$\begin{cases} D_t^\gamma x(t) = x(t), \\ D_t^\gamma y(t) = 2x^2(t), \\ D_t^\gamma z(t) = 3x(t)y(t), \end{cases} \quad 0 < \gamma < 1, \quad (4.5)$$

subject to initial conditions

$$x(0) = 1, \quad y(0) = 1, \quad z(0) = 0. \quad (4.6)$$

Operating Laplace transform on both sides of (17) and using (18), we get

$$\begin{cases} x(t) = 1 + \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L}[x(t)] \right], \\ y(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} \left[ \sum_{n=0}^{\infty} A_n(t) \right] \right], \\ z(t) = \mathcal{L}^{-1} \left[ \frac{1}{p^\gamma} \mathcal{L} \left[ \sum_{n=0}^{\infty} B_n(t) \right] \right], \end{cases} \quad (4.7)$$

where the nonlinear terms  $2x^2(t)$  and  $3x(t)y(t)$  in the system are expressed via Daftardar-Jafari polynomials  $A_n(t)$  and  $B_n(t)$ , respectively.

In view of LDJM,

$$\begin{aligned} x_0(t) &= 1, & y_0(t) &= 1, & z_0(t) &= 0, \\ x_1(t) &= \frac{t^\gamma}{\Gamma(\gamma+1)}, & y_1(t) &= \frac{2t^\gamma}{\Gamma(\gamma+1)}, & z_1(t) &= \frac{3t^\gamma}{\Gamma(\gamma+1)}, \\ x_2(t) &= \frac{t^{2\gamma}}{\Gamma(2\gamma+1)}, & y_2(t) &= \frac{4t^{2\gamma}}{\Gamma(2\gamma+1)} + \frac{2\Gamma(2\gamma+1)t^{3\gamma}}{\Gamma^2(\gamma+1)\Gamma(3\gamma+1)}, \\ z_2(t) &= \frac{9t^{2\gamma}}{\Gamma(2\gamma+1)} + \frac{6\Gamma(2\gamma+1)t^{3\gamma}}{\Gamma^2(\gamma+1)\Gamma(3\gamma+1)}, \\ x_3(t) &= \frac{t^{3\gamma}}{\Gamma(3\gamma+1)}, \\ y_3(t) &= \frac{4t^{3\gamma}}{\Gamma(3\gamma+1)} + \frac{4\Gamma(3\gamma+1)t^{4\gamma}}{\Gamma(\gamma+1)\Gamma(2\gamma+1)\Gamma(4\gamma+1)} + \frac{2\Gamma(4\gamma+1)t^{5\gamma}}{\Gamma^2(\gamma+1)\Gamma(5\gamma+1)}, \\ z_3(t) &= \frac{15t^{3\gamma}}{\Gamma(3\gamma+1)} + \frac{1}{\Gamma(\gamma+1)\Gamma(4\gamma+1)} \left( \frac{2\Gamma(3\gamma+1)}{\Gamma(\gamma+1)} + \frac{18\Gamma(3\gamma+1)}{\Gamma(2\gamma+1)} \right) t^{4\gamma} \\ &\quad + \frac{6\Gamma(2\gamma+1)\Gamma(4\gamma+1)t^{5\gamma}}{\Gamma^3(\gamma+1)\Gamma(3\gamma+1)\Gamma(5\gamma+1)}. \end{aligned}$$

Thus, the approximate solution of system (17) is

$$\begin{aligned} x(t) &= 1 + \frac{t^\gamma}{\Gamma(\gamma+1)} + \frac{t^{2\gamma}}{\Gamma(2\gamma+1)} + \frac{t^{3\gamma}}{\Gamma(3\gamma+1)} + \dots, \\ y(t) &= 1 + \frac{2t^\gamma}{\Gamma(\gamma+1)} + \frac{4t^{2\gamma}}{\Gamma(2\gamma+1)} + \left( 4 + \frac{2\Gamma(2\gamma+1)}{\Gamma^2(\gamma+1)} \right) \frac{t^{3\gamma}}{\Gamma(3\gamma+1)} \\ &\quad + \frac{4\Gamma(3\gamma+1)t^{4\gamma}}{\Gamma(\gamma+1)\Gamma(2\gamma+1)\Gamma(4\gamma+1)} + \frac{2\Gamma(4\gamma+1)t^{5\gamma}}{\Gamma^2(2\gamma+1)\Gamma(5\gamma+1)} + \dots, \\ z(t) &= \frac{3t^\gamma}{\Gamma(\gamma+1)} + \frac{9t^{2\gamma}}{\Gamma(2\gamma+1)} + \left( \frac{6\Gamma(2\gamma+1)}{\Gamma^2(\gamma+1)} + 15 \right) \frac{t^{3\gamma}}{\Gamma(3\gamma+1)} \\ &\quad + \frac{1}{\Gamma(\gamma+1)\Gamma(4\gamma+1)} \left( \frac{2\Gamma(3\gamma+1)}{\Gamma(\gamma+1)} + \frac{18\Gamma(3\gamma+1)}{\Gamma(2\gamma+1)} \right) t^{4\gamma} \\ &\quad + \frac{6\Gamma(2\gamma+1)\Gamma(4\gamma+1)t^{5\gamma}}{\Gamma^3(\gamma+1)\Gamma(3\gamma+1)\Gamma(5\gamma+1)} + \dots. \end{aligned}$$

In particular, if we put  $\gamma = 1$ , we have the exact solution in closed form

$$\begin{cases} x(t) = 1 + t + \frac{t^2}{2!} + \frac{t^3}{3!} + \dots \cong e^t, \\ y(t) = 1 + 2t + 2t^2 + \frac{4}{3}t^3 + \frac{1}{2}t^4 + \frac{2}{5}t^5 + \dots \cong e^{2t}, \\ z(t) = 3t + \frac{9}{2}t^2 + \frac{23}{6}t^3 + \frac{11}{4}t^4 + \frac{2}{5}t^5 + \dots \cong e^{3t} - 1. \end{cases} \tag{4.8}$$

Tabl .4 presents the numerical results for the approximate solution of Example .4.2 (4-term solution) at  $\gamma = 1$  and compares them with the exact solution and LRPS method [1]. Tab .5 shows a comparison of absolute errors. Furthermore, Tab .6 displays the approximate solutions obtained by LDJM, for Example 4.2 at different values of  $\gamma$ . Fig1b . illustrates the comparison between the LDJM and LRPS methods for  $y(t)$  and  $z(t)$  at  $\gamma = 1$ .

Table 4: A Comparison of numerical results obtained for Example 2 using LDJM and LRPS when  $\gamma = 1$ .

t	Exact solution			LDJM solution			LRPS solution		
	x(t)	y(t)	z(t)	x(t)	y(t)	z(t)	x(t)	y(t)	z(t)
0	1.000000	1.000000	0.000000	1.000000	1.000000	0.000000	1.000000	1.000000	0.000000
0.1	1.105171	1.221403	0.349859	1.105167	1.221387	0.349112	1.105167	1.221000	0.348833
0.2	1.221403	1.491825	0.822119	1.221333	1.491595	0.815195	1.221333	1.488000	0.810667
0.3	1.349859	1.822119	1.459603	1.349500	1.821022	1.431747	1.349500	1.807000	1.408500
0.4	1.491825	2.225541	2.320117	1.490667	2.222229	2.239829	1.490667	2.184000	2.165333
0.5	1.648721	2.718282	3.481689	1.645833	2.710417	3.288542	1.645833	2.625000	3.104167

Table 5: A Comparison of absolute errors obtained for Example 2 using LDJM and LRPS when  $\gamma = 1$ .

t	LDJM solution			LRPS solution		
	x(t)	y(t)	z(t)	x(t)	y(t)	z(t)
0	0.000000	0.000000	0.000000	0.000000	0.000000	0.000000
0.1	0.000004	0.000015	0.000080	0.000004	0.000403	0.001025
0.2	0.000069	0.000230	0.001591	0.000069	0.003825	0.011452
0.3	0.000359	0.001097	0.009856	0.000359	0.015119	0.051103
0.4	0.001158	0.003312	0.037621	0.001158	0.041541	0.154784
0.5	0.002888	0.007865	0.109814	0.002888	0.093282	0.377522

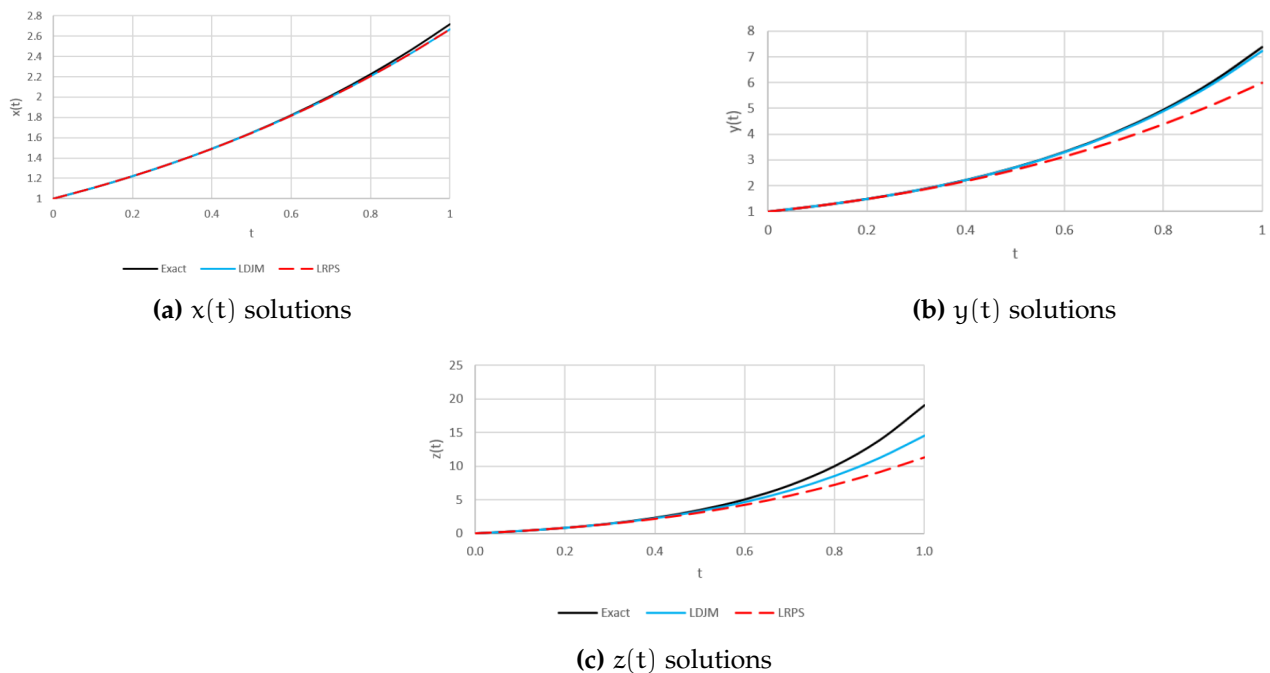


Figure 2: Comparison of  $x(t)$ ,  $y(t)$ , and  $z(t)$  solutions between exact, LDJM, and LRPS methods for Example 4.2

Table 6: Numerical results obtained for Example 2 using LDJM for different values of  $\gamma$ .

t	$\gamma = 0.9$			$\gamma = 0.7$			$\gamma = 0.5$		
	x(t)	y(t)	z(t)	x(t)	y(t)	z(t)	x(t)	y(t)	z(t)
0	1.000000	1.000000	0.000000	1.000000	1.000000	0.000000	1.000000	1.000000	0.000000
0.1	1.136481	1.284988	0.446971	1.225542	1.464073	0.717750	1.362201	1.736270	1.124769
0.2	1.267463	1.588855	0.976899	1.381811	1.823546	1.344304	1.527636	2.111597	1.774780
0.3	1.405977	1.948203	1.672894	1.533116	2.221817	2.136360	1.673194	2.496572	2.555300
0.4	1.555515	2.383411	2.599905	1.690047	2.699275	3.206975	1.817722	2.951133	3.620373
0.5	1.718374	2.915930	3.830516	1.858165	3.291506	4.673168	1.969900	3.521252	5.118808

**Comment:** Based on the above tables, it is evident that the proposed method exhibits a clear superiority in terms of result accuracy and convergence speed, particularly in nonlinear equations characterized by computational complexities, where the proposed method demonstrated stability and consistency in results.

### 5. Conclusion

This study has developed a novel analytical solution for nonlinear fractional autonomous dynamical systems using the Caputo fractional derivative. The proposed method, based on a combination of the Laplace transform and the Daftardar-Jafari method, has demonstrated high accuracy and efficiency in solving these systems. The results of this study confirm the effectiveness of the proposed method in solving similar problems and its potential applications in various scientific and engineering fields. Future research directions will focus on extending the application of this technique to more complex systems and exploring its potential in solving real-world problems.

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